Is there a natural state for Abelian Chern-Simons theory?

Simone Murro

Department of Mathematics University of Regensburg

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Motivations and goals of my talk

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"What is a QFT?" for a deeper understanding mathematical axioms for QFT
     Algebraic QFT :
               \checkmark \mathcal{M} \supset \mathcal{O} \mapsto \mathcal{A}(\mathcal{O}) \subset \mathcal{A}(\mathcal{M}) := \bigcup_{\mathcal{O}} \mathcal{A}(\mathcal{O})
                            - Isotony: if \mathcal{O} \subseteq \mathcal{O}' \Longrightarrow \mathcal{A}(\mathcal{O}) \subseteq \mathcal{A}(\mathcal{O}')
                            - Causality: if \mathcal{O} \cap J(\mathcal{O}') = \emptyset \Longrightarrow [\mathcal{A}(\mathcal{O}), \mathcal{A}(\mathcal{O}')] = 0
                            - Covariance: isometry \iota of M_1 \Longrightarrow \alpha_{\iota} \in \operatorname{Aut}(\mathcal{A}) s.t. \alpha_{\iota} \mathcal{A}(\mathcal{O}) = \mathcal{A}(\iota(\mathcal{O}))
               \checkmark \omega : \mathcal{A}(\mathcal{M}) \to \mathbb{C} s.t. normalised and positive \Rightarrow GNS-theorem (\mathcal{H}, \pi, \Omega)

    Locally covariant QFT

               \checkmark \mathfrak{A} : \mathsf{Loc} \to \mathsf{Alg}
                            - Locality: f: \mathcal{M} \to \mathcal{M}' \Longrightarrow f: \mathcal{A}(\mathcal{M}) \to \mathcal{A}(\mathcal{M}') injective
                           - Causality: \mathcal{M}_1 \xrightarrow{f_1} \mathcal{M} \xleftarrow{f_2} \mathcal{M}_2 caus. disjoint \Longrightarrow [f_1 \mathcal{A}(\mathcal{M}_1), f_2 \mathcal{A}(\mathcal{M}_2)] = 0
                            - Time-slice: f: \mathcal{M} \to \mathcal{M}' s.t. f(\mathcal{M}) \supset \Sigma' \Longrightarrow \mathfrak{f} is isomorphism
                     i natural state? f: \mathcal{M} \to \mathcal{M}' and \omega_{M}, \omega_{M'} invariant \omega_{M'} \circ f = \omega_{M}
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GOAL: Does a natural state exists for a Topological QFT?

Outline of the Talk

- Abelian Chern-Simons theory
- Quantization in the algebraic approach
- Invariant functionals on compact surfaces
- ¿ Natural states?

Based on:

► C. Dappiaggi, S.M., A. Schenkel, (arXiv:1612.04080)

The moduli space of flat U(1)-connections

- We consider $\mathcal{M} \simeq \mathbb{R} \times \Sigma$ without boundary and with dim $\Sigma = 2$
- The action of Abelian Chern–Simons theory

$$S = \frac{1}{4\pi} \int_{M} \operatorname{tr}(A \wedge dA) \quad \Rightarrow \quad 0 = \frac{\delta S}{\delta A} = \frac{1}{2\pi} dA$$

• The moduli space of flat U(1)-connection:

$$\mathsf{Flat}_{\mathit{U}(1)} := \frac{\Omega^1_{\mathit{d}}(\mathit{M})}{\Omega^1_{\mathbb{Z}}(\mathit{M})} \simeq \frac{\mathit{H}^1(\mathit{M};\mathbb{R})}{\mathit{H}^1(\mathit{M};\mathbb{Z})} \simeq \frac{\mathit{H}^1(\Sigma;\mathbb{R})}{\mathit{H}^1(\Sigma;\mathbb{Z})} \simeq \frac{\Omega^1_{\mathit{d}}(\Sigma)}{\Omega^1_{\mathbb{Z}}(\Sigma)}$$

• In the categorical language, the assignment of the moduli spaces is a functor:

$$\begin{split} \mathsf{Flat}_{\mathit{U}(1)} : \ \mathsf{Man}_{2}^{\mathsf{op}} &\to \mathsf{Ab} \\ \mathsf{Flat}_{\mathit{U}(1)}(f) := f^{*} : \frac{\Omega_{d}^{1}(\Sigma')}{\Omega_{\mathbb{Z}}^{1}(\Sigma')} &\to \frac{\Omega_{d}^{1}(\Sigma)}{\Omega_{\mathbb{Z}}^{1}(\Sigma)} \qquad [\mathsf{A}'] \mapsto [f^{*}\mathsf{A}'] \end{split}$$

- Obj(Man₂)={2-dimensional oriented manifolds}
- Hom(Man₂)={orientation preserving open embeddings}

Observable for Abelian Chern-Simons theory

• As basic osservables we take all group characters := $\operatorname{Hom}(\operatorname{Flat}_{U(1)}(\Sigma), U(1))$:

given any
$$\varphi \in \Omega^1_c(\Sigma)$$
 $A \mapsto \exp\left(2\pi\imath \int_{\Sigma} \varphi \wedge A\right)$

- This character descends to the quotient if and only if $\int_{\Sigma} \varphi \wedge \Omega^1_{\mathbb{Z}}(\Sigma) \subseteq \mathbb{Z}$
- Since $d\Omega^0(\Sigma) \subseteq \Omega^1_{\mathbb{Z}}(\Sigma)$ is a subgroup, Stokes' lemma implies that any $\varphi \in \Omega^1_{c,d}(\Sigma)$
- Because each exact $\varphi = d\chi \in d\Omega^0_c(\Sigma)$ yields a trivial group character, then

$$\text{group character} \simeq \textit{H}^1_c(\Sigma;\mathbb{Z}) := \left\{ [\varphi] \in \textit{H}^1_c(\Sigma;\mathbb{R}) : \int_{\Sigma} \varphi \wedge \textit{H}^1(\Sigma) \subseteq \mathbb{Z} \right\}$$

• The assignment of the character groups is a functor

$$\begin{split} \mathcal{O} := \left(\mathcal{H}^1_c(-;\mathbb{Z}), \tau \right) : \mathsf{Man}_2 \to \mathsf{PAb} \\ \mathcal{H}^1_c(f;\mathbb{Z}) := f_* \, : \, \mathcal{H}^1_c(\Sigma;\mathbb{Z}) \, \longrightarrow \, \mathcal{H}^1_c(\Sigma';\mathbb{Z}) \, \, , \quad [\varphi] \longmapsto [f_*\varphi] \\ \tau_\Sigma : \, \mathcal{H}^1_c(\Sigma;\mathbb{Z}) \times \mathcal{H}^1_c(\Sigma;\mathbb{Z}) \longrightarrow \mathbb{R} \qquad \tau([\varphi], [\widetilde{\varphi}])_\Sigma = \int_\Sigma \varphi \wedge \widetilde{\varphi} \\ \tau_{\Sigma'} \big(f_*[\varphi], f_*[\widetilde{\varphi}] \big) = \int_{\Sigma'} (f_*\varphi) \wedge (f_*\widetilde{\varphi}) = \int_\Sigma \varphi \wedge (f^*f_*\widetilde{\varphi}) = \int_\Sigma \varphi \wedge \widetilde{\varphi} = \tau_\Sigma \big([\varphi], [\widetilde{\varphi}] \big) \, \, , \end{split}$$

Quantization of Abelian Chern-Simons theory

ullet Quantization is achieved composing $\mathcal{O}:\mathsf{Man}_2\to\mathsf{PAb}$ with $\mathfrak{CCR}:\mathsf{PAb}\to\mathsf{CAlg}$

$$\mathcal{A} := \mathfrak{CCR} \circ \mathcal{O} : \mathsf{Man}_2 \to \mathsf{CAlg}$$

• We construct a *-algebra $\Delta:=$ span $\{W_{[\varphi]}\,|\, [\varphi]\in H^1_c(\Sigma;\mathbb{Z})\}$ where $W_{[\,\cdot\,]}$ satisfy

$$W_{[\varphi]} W_{[\widetilde{\varphi}]} := e^{-i\hbar \, au_{\Sigma}([\varphi], [\widetilde{\varphi}])} W_{[\varphi]+[\widetilde{\varphi}]} \quad , \qquad W_{[\varphi]}^* := W_{-[\varphi]} \quad .$$

• We obtain a C^* -algebra taking the completion of Δ with respect to the norm

$$||a||^{m.r.} := \sup_{\omega \in \mathcal{F}} \sqrt{\omega(a^*a)}$$

where $\omega: \Delta \to \mathbb{C}$ is a state, namely $\omega(1_{\Delta}) = 1$ and $\omega(a^*a) \geq 0$.

- ullet We assume that $\hbar
 ot\in 2\pi \mathbb{Z}$ to avoid commutative C^* -algebras
- ullet The C^* -algebra homomorphism $\mathcal{A}(f):\mathcal{A}(\Sigma) o\mathcal{A}(\Sigma')$ is specified by

$$\mathcal{A}(f)\big(W_{[\,\cdot\,]}\big):=W'_{f_*[\,\cdot\,]}\ ,$$

where by W' we denote the Weyl symbols in $\mathcal{A}(\Sigma')$.

Invariant functionals on compact surfaces

- Any object Σ in Man₂ comes together with its automorphism Diff⁺(Σ) of Σ
- Because $H^1_c(\Sigma; \mathbb{Z})$ is discrete, $\mathsf{Diff}^+_0(\Sigma) \subseteq \mathsf{Diff}^+(\Sigma)$ is represented trivially

$$\mathsf{MCG}(\Sigma) := rac{\mathsf{Diff}^+(\Sigma)}{\mathsf{Diff}^+_0(\Sigma)} o \mathsf{Aut}ig(\mathcal{A}(\Sigma)ig) \qquad f \mapsto \mathcal{A}(f)$$

ullet For compact Σ , there exists a short exact sequence of groups

$$1 \ \longrightarrow \ \mathsf{Tor}(\Sigma) \ \longrightarrow \ \mathsf{MCG}(\Sigma) \ \longrightarrow \ \mathsf{Sp}(H^1(\Sigma;\mathbb{Z}),\tau_\Sigma) \ \longrightarrow 1$$

• The representation of the $\mathsf{MCG}(\Sigma)$ descends to a representation of $\mathsf{Sp}(H^1(\Sigma;\mathbb{Z}),\tau_\Sigma)$

$$\mathsf{Sp}(H^1(\Sigma;\mathbb{Z}), au_\Sigma)\longrightarrow \mathsf{Aut}(\mathcal{A}(\Sigma))\;,\;\; T\longmapsto \kappa_T:W_{[\varphi]}\mapsto W_{T[\varphi]}$$

• An invariant functional under the action of the symplectic group

$$\omegaig(W_{\mathcal{T}[arphi]}ig) = \omegaig(W_{[arphi]}ig) := egin{cases} 1 & ext{if } [arphi] = 0 \ K_{[arphi]} & ext{else} \end{cases}$$

• Since $Sp(H^1(\Sigma; \mathbb{Z}), \tau_{\Sigma}) \not\subset U(H^1(\Sigma; \mathbb{Z}), \mu)$, does not exist invariant Gaussian state

$$\omega(W_{[\varphi]}) = e^{-\mu([\varphi],[\varphi])}$$

Non-existence of natural states

No-go theorem: There exists *no natural state* or the functor $\mathcal{A}:\mathsf{Man}_2\to\mathsf{CAlg},$ namely a state for each Σ such that for all Man_2 -morphisms $f:\Sigma\to\Sigma'$ holds true:

$$\omega_{\Sigma'} \circ \mathcal{A}(f) = \omega_{\Sigma}$$

Sketch of the proof

- ullet Let us assume that there exists a natural state $\{\omega_\Sigma\}_{\Sigma\in\mathsf{Man}}$
- Consider the Man₂-diagram:

$$\mathbb{S}^2 \xleftarrow{f_1} \mathbb{R} \times \mathbb{T} \xrightarrow{f_2} \mathbb{T}^2$$

• The naturality of the state implies:

$$\omega_{\mathbb{S}^2} \circ \mathcal{A}(\mathit{f}_1) = \omega_{\mathbb{R} imes \mathbb{T}} = \omega_{\mathbb{T}^2} \circ \mathcal{A}(\mathit{f}_2)$$

- Because of $H^1_c(\mathbb{S}^2;\mathbb{Z})=0$, then $\mathcal{A}(\mathbb{S}^2)\simeq\mathbb{C}$ and hence $\omega_{\mathbb{S}^2}=\mathrm{id}_{\mathbb{C}}$ is unique on \mathbb{C}
- $\bullet \ \ \text{We can choose} \ \textit{f}_2 \ \text{such that} \ \ \textit{W}_n^{\mathbb{R} \times \mathbb{T}} \mapsto \textit{W}_{(n,0)}^{\mathbb{T}^2} \ \text{we obtain that} \ \ \omega_{\mathbb{T}^2}\big(\textit{W}_{(n,0)}^{\mathbb{T}^2}\big) = 1$
- Choosing $a=\alpha_1\,1+\alpha_2\,W_{(1,1)}^{\mathbb{T}^2}+\alpha_3\,W_{(0,1)}^{\mathbb{T}^2}\in\mathcal{A}(\mathbb{T}^2)$ the functional $\omega_{\mathbb{T}^2}(a^*a)<0$

Q.E.D